



# CHARACTERIZATION OF THE NATURALLY FRACTURED RESERVOIR PARAMETERS IN INFINITE-CONDUCTIVITY HYDRAULICALLY-FRACTURED VERTICAL WELLS BY TRANSIENT PRESSURE ANALYSIS

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## ABSTRACT

It has become common to hydraulically fracture a naturally fractured formation to increase the well's production potential. Since the mass transfer between fractured network and hydraulic fracture is much higher than that from matrix to fractures, the hydraulic fracture-fracture network interporosity flow parameter is much higher than that of matrix-fracture network. As a result, the transition period behavior from naturally fractured to homogeneous takes place before radial flow regime during the early bilinear, linear or elliptical (birradial) flow regimes. The purpose of this paper is to provide expressions by both conventional analysis and *TDS* technique for characterizing the naturally fractured parameters when the transition period interrupts the response of an infinite-conductivity fracture. The developed expressions for both methodologies were satisfactorily tested with simulated examples.

**Keywords:** pressure transient analysis, interporosity flow parameter, *TDS* technique, elliptical flow, naturally fractured reservoirs.

## 1. INTRODUCTION

Even wells in naturally fractured formations are being subjected to hydraulic fracturing since the bulk permeability is low and well productivity is not as high as expected. Typical such cases have been seen in some wells of the Orinoco basin foothill in Colombia where the wells were fractured just after being drilled.

Tiab (1994) was the first to identify the elliptical flow regime in pressure test data of infinite-conductivity hydraulically fractured vertical wells. He called it "birradial flow", provided the governing pressure and pressure derivative equations and implemented the *TDS* technique for its characterization. This flow regime can also be seen in horizontal wells before the late radial flow regime shows up and is due mainly to horizontal anisotropy. The last model for the elliptical flow in horizontal wells was presented by Martinez, Escobar and Bonilla (2012).

The original pressure governing model for elliptical flow introduced by Tiab (1994) involves the drainage area, which may affect the calculations if the pressure test is too short for the development of the late pseudosteady-state period. An experienced user of the *TDS* technique may deal with that problem using the point of intersection between the radial flow and the birradial flow regime lines, Tiab (1994). However, to overcome this problem, Escobar, Bonilla and Ghisays-Ruiz (2014) formulated a new elliptical/birradial flow model for vertical wells in either fractured or unfractured formations.

Although Tiab's model, Tiab (1994), works perfectly, it requires the knowledge of well-drainage area for estimating the half-fracture length. However, sometimes this condition cannot be met because the pressure test needs to be run long enough for the development of the late time pseudosteady-state period for the determination of the drainage area. The new model presented by Escobar *et al.* (2014) excludes the reservoir drainage area and slightly modifies Tiab's model, Tiab (1994), to account for naturally-fractured double-porosity systems.

Tiab and Bettam (2007) presented a practical interpretation of the pressure behavior of a finite-conductivity hydraulically fractured vertical well located in a naturally fractured reservoir. The interpretation is based on analytical equations derived to determine permeability, fracture storage capacity ratio, interporosity flow coefficient, skin and wellbore storage from the pressure derivative plot without using type-curve matching technique. In other words, they implemented the *TDS* technique, Tiab (1993), for such systems. Part of the work presented by Escobar, Martinez and Montealegre (2009) was focused on the implementation of conventional analysis for the work presented by Tiab and Bettam (2007).

In their work, Tiab and Bettam (2007) developed expressions for bilinear and linear flow regimes but excluded the presence of the elliptical or birradial flow regime. Thereafter, this work constitutes an extension of



the research of Tiab and Bettam (2007) for infinite-conductivity fractures where early birradial flow is seen. This implementation was satisfactorily tested with synthetic examples.

## 2. MATHEMATICAL FORMULATION

### 2.1. Governing equation

The dimensionless pressure behavior in the Laplace domain for a hydraulically fractured well in double-porosity reservoirs was introduced by Ozkan and Raghavan (1988):

$$\bar{P}_D(s) = \frac{1}{2s\sqrt{sf(s)}} \int_0^{s\sqrt{sf(s)}[1 \pm 0.732]} K_0(\alpha) d\alpha \quad (1)$$

For a pseudosteady-state matrix flow model, the function of storage capacity and interporosity flow parameter is given by:

$$f(s) = \omega + \frac{(1-\omega)\lambda_f}{(1-\omega)s + \lambda_f} \quad (2)$$

Escobar *et al.* (2014) provided the following pressure and pressure derivative governing equations for elliptical flow regime:

$$P_D = \frac{25}{9} \left( \frac{\pi t_{Dxf}}{26\xi} \right)^{0.36} \quad (3)$$

$$t_D * P_D' = \left( \frac{\pi t_{Dxf}}{26\xi} \right)^{0.36} \quad (4)$$

When  $\xi = 1$ , Equations (3) and (4) account for homogeneous reservoirs. For the case of naturally-fractured formations,  $\xi = \omega$ , which is the dimensionless storativity coefficient. Consequently, for the purpose of this work,  $\xi$  will be replaced by the storativity ratio.

### 2.2. Dimensionless quantities

The dimensionless time, based on half-fracture length, is given below:

$$t_{Dxf} = \frac{0.000263kt}{\phi\mu c_t x_f^2} \quad (5)$$

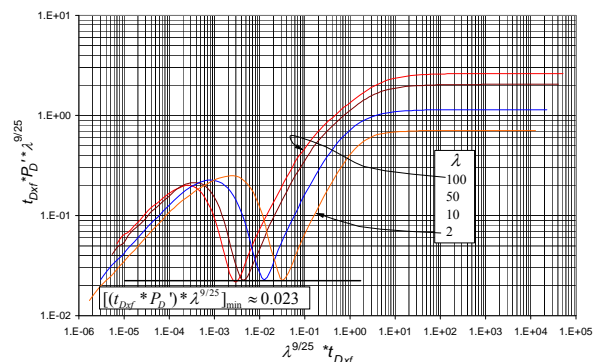
The dimensionless pressure and pressure derivative parameters for oil reservoirs are given by:

$$P_D = \frac{kh\Delta P}{141.2q\mu B} \quad (6)$$

$$t_D * P_D' = \frac{kh(t * \Delta P')}{141.2q\mu B} \quad (7)$$

Finally, the dimensionless fracture conductivity introduced by Cinco-Ley, Samaniego and Dominguez (1976) is defined as:

$$C_{fD} = \frac{k_f w_f}{k x_f} \quad (8)$$



**Figure-1.** Effect of the interporosity flow parameter on the pressure and pressure derivative during the elliptical flow regime,  $\omega = 0.01$ .

### 2.3. TDS technique

By plugging the dimensionless parameters in Equations (3) and (4) and solving for the half-fracture length, we obtain:

$$x_f = 22.5632 \left( \frac{qB}{h\Delta P_{BR1}} \right)^{1.3889} \sqrt{\frac{1}{\omega\phi c_t} \left( \frac{\mu}{k} \right)^{1.778}} \quad (9)$$

$$x_f = 5.4595 \left( \frac{qB}{h(t * \Delta P')_{BR1}} \right)^{1.3889} \sqrt{\frac{1}{\omega\phi c_t} \left( \frac{\mu}{k} \right)^{1.778}} \quad (10)$$

Figures-1 through 4 were generated using the analytical solution provided by Ozkan and Raghavan (1988), Equation (1), which considers flow between hydraulic fracture and fracture network.

As seen in Figure-1, once the transition period vanishes, the reservoir behaves as if it were homogeneous. In this case, Equations (9) and (10) become:



$$x_f = 22.5632 \left( \frac{qB}{h\Delta P_{2BR1}} \right)^{1.3889} \sqrt{\frac{1}{\phi c_i} \left( \frac{\mu}{k} \right)^{1.778}} \quad (11)$$

$$x_f = 2.032 \left( \frac{qB}{(t^* \Delta P')_{L1} h \sqrt{\omega}} \right) \sqrt{\frac{\mu}{(\phi c_i) k}} \quad (14)$$

$$x_f = 5.4595 \left( \frac{qB}{h(t^* \Delta P')_{2BR1}} \right)^{1.3889} \sqrt{\frac{1}{\phi c_i} \left( \frac{\mu}{k} \right)^{1.778}} \quad (12)$$

Fracture conductivity can be found using an expression presented by Tiab (2003);

$$k_f w_f = \frac{3.31739k}{e^s - \frac{1.92173}{x_f}} \quad (15)$$

Where 2BR1 corresponds to the straight line drawn on the second elliptical flow regime which occurs after the transition period. However, before the transition period either elliptical flow or linear flow may take place. If linear flow occurs at early time, then Equations (9) and (10) must be replaced by these two expressions developed by Tiab and Bettam (2007);

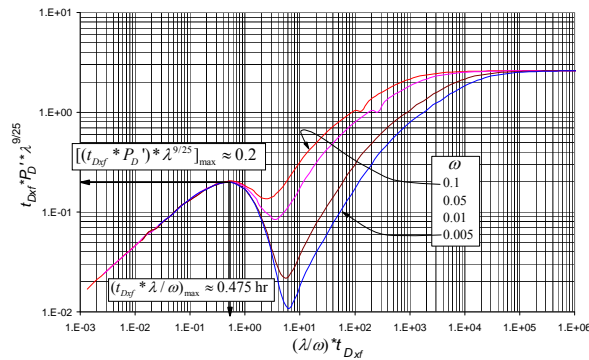
It is observed in Figure-1 that the minimum point of the pressure derivative times the interporosity flow parameter to the power 0.36 is a constant:

$$x_f = 4.064 \left( \frac{qB}{\Delta P_{L1} h \sqrt{\omega}} \right) \sqrt{\frac{\mu}{(\phi c_i) k}} \quad (13)$$

$$(t_{Dxf}^* P_D')_{\min} \lambda_f^{9/25} \approx 0.023 \quad (16)$$

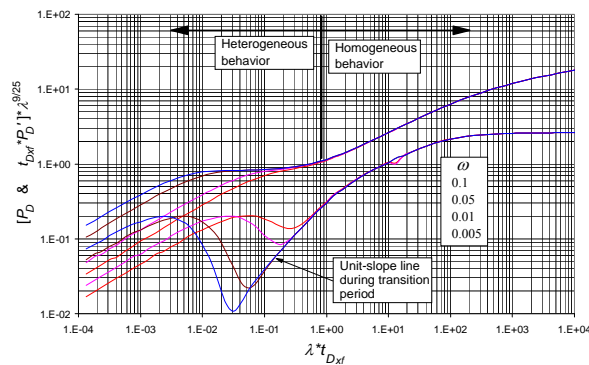
The following expression to determine the interporosity flow parameter is obtained once Equation (7) is substituted in Equation (16):

$$\lambda_f \approx 26.364 \left[ \frac{q\mu B}{kh(t^* \Delta P')_{\min}} \right]^{25/9} \quad (17)$$



**Figure-2.** Correlation of the maximum and minimum pressure derivative coordinates during elliptical flow period,  $\lambda_f = 100$ .

Figure-2 contains a log-log plot of pressure and pressure derivative both multiplied by the interporosity flow parameter raised to the power 9/25 (or 0.36) versus dimensionless time multiplied by interporosity flow parameter and divided by the storativity ratio. During the early heterogeneous behavior, all curves unify no matter the values of storativity ratio. A maximum point is displayed before the transition period which coordinates are given by:



**Figure-3.** Effect of the storativity ratio on the pressure and pressure derivative during the elliptical flow regime,  $\lambda_f = 100$ .

$$[(t_{Dxf}^* P_D') \lambda_f^{9/25}]_{\max} \approx 0.2 \quad (18)$$

$$\left( \frac{t_{Dxf} \cdot \lambda_f}{\omega} \right)_{\max} \approx 0.475 \quad (19)$$

Once Equation (7) is substituted into Equation (18), we obtain:

$$\lambda_f \approx 10719.63 \left[ \frac{q\mu B}{kh(t^* \Delta P')_{\max}} \right]^{25/9} \quad (20)$$

Substituting Equation (5) into Equation (19) yields two expressions to find either interporosity flow parameter or storativity ratio one as a function of the other so;



$$\omega = \frac{\lambda_f k t_{\max}}{1801.29 \phi \mu c_i x_f^2} \quad (21)$$

$$\lambda_f = \frac{1801.29 \omega \phi \mu c_i x_f^2}{k t_{\max}} \quad (22)$$

Figure-3 was prepared to unify the unit-slope line showing up during the transition period. Such straight line was fitted to:

$$(t_{Dxf} * P_D') = 0.3862 t_{Dxf} \lambda_f^{16/25} \quad (23)$$

After substituting the dimensionless parameter and solving for the interporosity flow parameter, it yields:

$$\lambda_f = 755.94 \left( \frac{\phi c_i x_f^2 h (t^* \Delta P')_{us}}{q B t_{us}} \right)^{25/16} \quad (24)$$

Equation (24) works using the coordinates of any point along the unit-slope line. During radial flow regime, the value of the dimensionless pressure derivative is one half. Equating this value to Equation (23) and solving for the interporosity flow parameter, we get:

$$\lambda_f = 585120.28 \left( \frac{\phi \mu c_i x_f^2}{k t_{usi}} \right)^{25/16} \quad (25)$$

Where  $t_{usi}$  is the intersection point of the unit-slope line and the radial flow regime.

From observation of the minimum point in Figure-2, an expression with a correlation coefficient of 0.999 was obtained:

$$\omega = 0.0002 \left( \frac{t_{\min}}{t_{\max}} \right)^2 - 0.0155 \left( \frac{t_{\min}}{t_{\max}} \right) + 0.2487 \quad (26)$$

Dividing Equation (10) by Equation (12) and solving for the storativity ratio, we get:

$$\omega = \left[ \frac{(t^* \Delta P')_{2BR1}}{(t^* \Delta P')_{BR1}} \right]^{2.7778} \quad (27)$$

Equation (27) is valid whenever birradial/elliptical flow exists at both side of the transition period. However, in infinite-conductivity fracture, it is possible that the linear flow occurs at early time, and then

is interrupted by the transition period, followed by the elliptical or birradial flow. In such cases, dividing Equation (14) by Equation (12) leads to:

$$\omega = 0.13853 \left( \frac{h}{qB} \right)^{0.778} \left[ \frac{(t^* \Delta P')_{2BR1}^{1.3889}}{(t^* \Delta P')_{L1}} \right]^2 \left( \frac{k}{\mu} \right)^{0.778} \quad (28)$$

The point of intersection between the early birradial line, Equation (4), and the unit-slope line during the transition period, Equation (23), leads to the following expression:

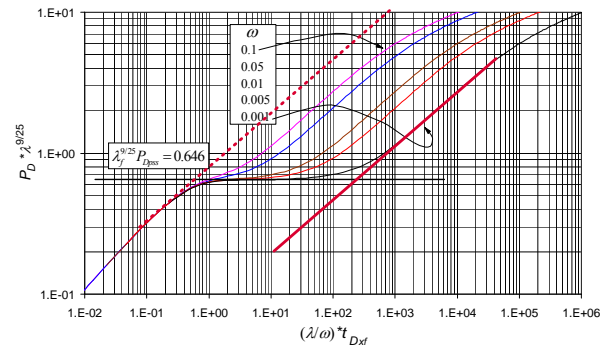
$$\omega = \frac{1}{1456.69} \left( \frac{k \lambda_f t_{usBR1}}{\phi \mu c_i x_f^2} \right)^{9/16} \quad (29)$$

Tiab and Bettam (2007) introduced the governing pressure derivative equation for early linear flow:

$$t_D * P_D' = \frac{1}{2} \sqrt{\frac{t_{Dxf}}{\omega}} \quad (30)$$

Which, in combination with Equation (23), leads to:

$$\omega = \frac{6356.314 \phi \mu c_i x_f^2}{k \lambda_f^{32/25} t_{usLi}} \quad (31)$$



**Figure-4.** Dimensionless pressure times the interporosity flow parameter to the power 0.36 versus dimensionless time multiplied by  $\lambda_f$ ,  $\lambda_f = 100$ .

#### 2.4. Conventional analysis

As seen in Figure-4, there is a point in which the dimensionless pressure converges to a value during the transition period. From there we can read that:

$$\lambda_f^{9/25} P_{Dps} = 0.833 \quad (32)$$



Substituting Equation (32) in and solving for the interporosity flow parameter yields:

$$\lambda_f = 564069.34 \left[ \frac{q\mu B}{kh\Delta P_{pss}} \right]^{25/9} \quad (33)$$

As seen in Figure-4, once the transition period vanishes a line with slope of 0.4 is developed. Draw such line and draw a parallel line throughout the points just before the development of the transition period. Read two points in both lines at the same time value and take its ratio, which is used in the following fitting equation that has a correlation coefficient of 0.999. The storativity ratio is found from such separation:

$$\omega = 0.3625 e^{-0.9877 \frac{\Delta P_{high}}{\Delta P_{low}}} \quad (34)$$

### 3. SYNTHETIC EXAMPLES

#### 3.1. Synthetic example 1

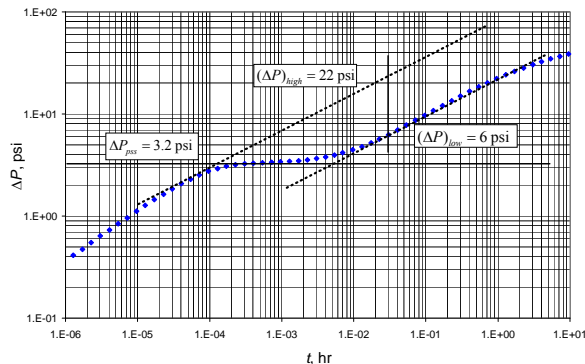
Figure-5 reports simulated pressure drop versus time data obtained with the information given on the second column of Table-1. It is required to find the naturally-fractured reservoir parameters.

Solution by Conventional Analysis

In Figure-5 a horizontal line was drawn going through the transition points. The following value of pressure drop was read:

$$\Delta P_{pss} = 3.2 \text{ psi}$$

Using the above value in Equation (33), it yields:



**Figure-5.** Pressure drop versus time log-log plot for example 1.

**Table-1.** Input data for examples.

Parameter	Examples	
	1	2
$k$ , md	5	1
$h$ , ft	100	20
$\phi$ , %	5	10
$B$ , rb/STB	1.2	1.1
$\mu$ , cp	1.5	3
$c_t$ , 1/psi	$1 \times 10^{-4}$	$1 \times 10^{-4}$
$q$ , BPD	30	50
$x_f$ , ft	10	30
$\lambda_f$	50	100
$\omega$	0.01	0.005

$$\lambda_f = 564069.34 \left[ \frac{(50)(1.5)(1.2)}{(5)(100)(3.2)} \right]^{25/9} = 46.05$$

Then, a 0.4-slope line is drawn after the transition period. Another parallel line goes just after the initiation of the transition period so the amplitude can be estimated. Two pressure drop values are read at the same time from both straight lines, so:

$$\Delta P_{high} = 22 \text{ psi}$$

$$\Delta P_{low} = 6 \text{ psi}$$

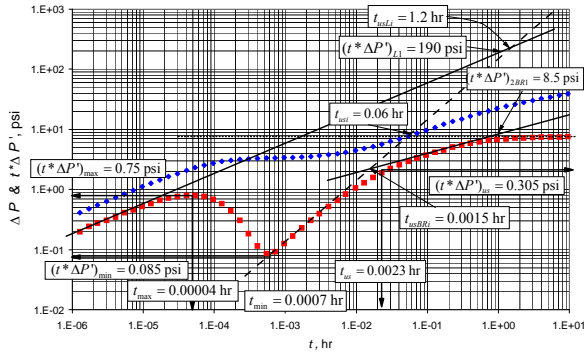
Equation (34) will lead to the estimation of the storativity ratio:

$$\omega = 0.3625 e^{-0.9877 \frac{22}{6}} = 0.0097 \approx 0.01$$

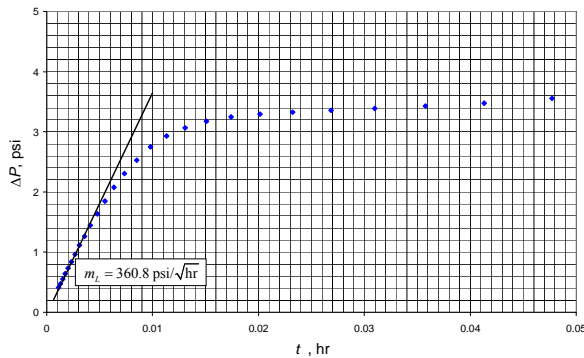
The half-fracture length is found from the slope of a Cartesian plot ( $m_L = 360.8 \text{ psi/hr}^{0.5}$ ) of pressure versus the square root of time, Figure-7, using Equation (A.7):

$$x_f = 4.064 \left( \frac{30(1.2)}{(100)(360.8)} \right) \sqrt{\frac{1.5}{(0.05)(0.0001)(5)(0.01)}} = 9.93 \text{ ft}$$

Solution by TDS Technique



**Figure-6.** Pressure and pressure derivative versus time log-log plot for example 1.



**Figure-7.** Cartesian plot of pressure versus the square root of time for example 1.

The following information was read from the pressure and pressure derivative log-log plot provided in Figure-6.

- |                                     |                                    |
|-------------------------------------|------------------------------------|
| $t_{max} = 0.00004$ hr              | $(t^* \Delta P')_{2BR1} = 8.5$ psi |
| $(t^* \Delta P')_{max} = 0.75$ psi  | $(t^* \Delta P')_{L1} = 190$ psi   |
| $(t^* \Delta P')_{min} = 0.085$ psi | $t_{us} = 0.0023$ hr               |
| $(t^* \Delta P')_{us} = 0.305$ psi  | $t_{usi} = 0.06$ hr                |
| $t_{usBR1} = 0.015$ hr              | $t_{usLi} = 1.2$ hr                |
| $(t^* \Delta P')_{2BR1} = 8.5$ psi  | $t_{min} = 0.0007$ hr              |

Using Equation (14);

$$x_f = 2.032 \left( \frac{30(1.2)}{(190)(100)\sqrt{0.01}} \right) \sqrt{\frac{1.5}{0.05(0.0001)(5)}} = 9.43 \text{ ft}$$

The half-fracture length is also found with Equation (12);

$$x_f = 5.4595 \left( \frac{30(1.2)}{100(8.5)} \right)^{1.3889} \sqrt{\frac{1}{0.05(0.0001)} \left( \frac{1.5}{5} \right)^{1.778}} = 10.37 \text{ ft}$$

The interporosity flow parameter is estimated with Equations (17), (20), (22), (24) and (25):

$$\lambda_f \approx 26.364 \left[ \frac{30(1.5)(1.2)}{5(100)(0.085)} \right]^{25/9} = 51.28$$

$$\lambda_f \approx 10719.63 \left[ \frac{30(1.5)(1.2)}{5(100)(0.75)} \right]^{25/9} = 49.24$$

$$\lambda_f = \frac{1801.29(0.1)(0.05)(1.5)(0.0001)(10^2)}{5(0.00004)} = 45.03$$

$$\lambda_f = 755.94 \left( \frac{(0.05)(0.0001)(10^2)(100)(0.305)}{(30)(1.2)(0.0023)} \right)^{25/16} = 54.6$$

$$\lambda_f = 585120.28 \left( \frac{(0.05)(0.0001)(10^2)}{5(0.06)} \right)^{25/16} = 50.3$$

The storativity ratio is estimated with Equations (26), (28), (29) and (31):

$$\omega = 0.0002 \left( \frac{0.0007}{0.00004} \right)^2 - 0.0155 \left( \frac{0.0007}{0.00004} \right) + 0.2487 = 0.039$$

$$\omega = 0.13853 \left( \frac{100}{30(1.2)} \right)^{0.778} \left[ \frac{(8.5)^{1.3889}}{(190)} \right]^2 \left( \frac{5}{1.5} \right)^{0.778} = 0.0083$$

$$\omega = \frac{1}{1456.69} \left( \frac{5(50)(0.015)}{(0.05)(1.5)(0.0001)(10^2)} \right)^{9/16} = 0.082$$

$$\omega = \frac{6356.314(0.05)(1.5)(0.0001)(10^2)}{(5)(50^{32/25})(1.2)} = 0.005$$

### 3.2. Synthetic example 2

Figure-8 presents synthetic pressure drop versus time data obtained with the information given on the third column of Table-1. It is required to find the naturally-fractured reservoir parameters.

Solution by Conventional Analysis

In Figure-8 a horizontal line was drawn going through the transition points. The following value of pressure drop was read:

- $\Delta P_{pss} = 184$  psi
- $\Delta P_{high} = 1600$  psi
- $\Delta P_{low} = 373$  psi



$m_L = 2465.84 \text{ psi/hr}^{0.5}$  read from Figure-10.

Equation (33) leads to an estimation of a value 101.36 for the interporosity flow parameter and Equation (34) provides a storativity ratio of 0.00523. A half-fracture length of 35.1 ft is found with Equation (A.7)

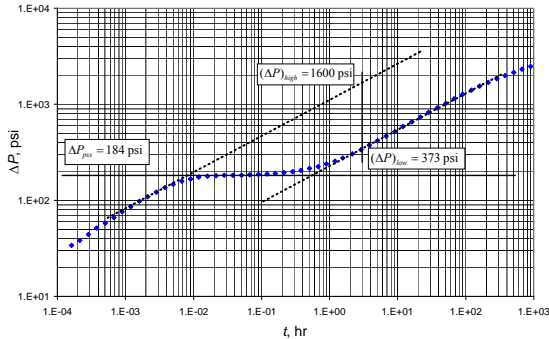


Figure-8. Pressure drop versus time log-log plot for example 1.

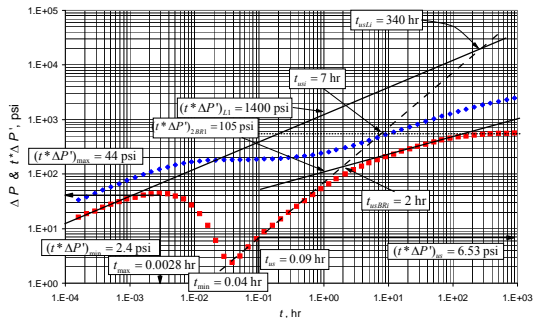


Figure-9. Pressure and pressure derivative versus time log-log plot for example 2.

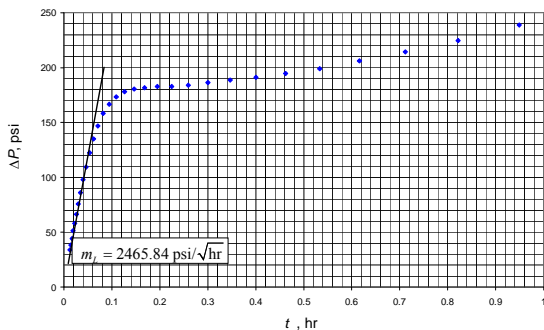


Figure-10. Cartesian plot of pressure versus the square root of time for example 2.

Solution by TDS Technique

The following information was read from the pressure and pressure derivative log-log plot provided in Figure-9.

- $t_{max} = 0.0028 \text{ hr}$
- $(t^* \Delta P^*)_{max} = 44 \text{ psi}$
- $(t^* \Delta P^*)_{min} = 2.4 \text{ psi}$
- $t_{usBR1} = 2 \text{ hr}$
- $(t^* \Delta P^*)_{us} = 6.53 \text{ psi}$
- $(t^* \Delta P^*)_{2BR1} = 105 \text{ psi}$
- $(t^* \Delta P^*)_{L1} = 1400 \text{ psi}$
- $t_{us} = 0.09 \text{ hr}$
- $t_{usL1} = 340 \text{ hr}$
- $t_{min} = 0.04 \text{ hr}$

The calculations were performed in the same way as in example 1 and a summary of the results is given in Table-3.

Table-2. Results from conventional analysis.

Parameter	Example 1	Input	Example 2	Input	Equation Number
$x_f$ , ft	9.93	10	35.1	30	A.7
$\lambda_f$	46.05	50	101.36	100	33
$\omega$	0.0097	0.01	0.00523	0.005	34

Table-3. Results from TDS technique.

Parameter	Example 1	Example 2	Equation Number
$x_f$ , ft	10	30	Input
	9.43	30.91	14
	10.37	29.12	12
$\lambda_f$	50	100	Input
	51.28	813.9	17
	49.24	102.5	20
	45.03	28.95	22
	54.6	79.96	24
$\omega$	50.3	99.03	25
	0.01	0.005	Input
	0.039	0.068	26
	0.083	0.0056	28
	0.082	0.069	29
	0.005	0.0034	31

4. COMMENTS ON THE RESULTS

The provided examples in this work demonstrate the applicability of the developed mathematical expressions. It is worth to remind that the estimation of the naturally fractured parameters is very sensitive, and as a result, deviation errors of one order of magnitude are allowed. The method provides satisfactory results especially for the interporosity flow parameter. If further



accuracy is required, then, use of computer program is highly recommended. These examples were performed by hand and the values from the plots were read by eye, consequently, deviation from input values is expected.

Equations for gas flow are given in Appendix B and the equations for the estimation of half-fracture length are given in Appendix A.

## CONCLUSIONS

- New mathematical expressions for both *TDS* technique and conventional analysis were introduced and successfully tested with synthetic cases for the full characterization of pressure tests conducted in naturally-fractured reservoirs whose wells have been subjected to hydraulic fracturing.
- For the case of the TDs technique, five expressions for the estimation of the interporosity flow parameter were developed. Four equations for the estimation of the storativity ratio were also introduced in this work.

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## Nomenclature

$A$	Draining area, ft <sup>2</sup>
$B$	Oil volume factor, rb/STB
$b_x$	Shortest distance from a lateral boundary to a well, ft
$C_{fD}$	Dimensionless fracture conductivity
$c_i$	Compressibility, 1/psi
$h$	Formation thickness, ft
$k$	Formation compressibility, md
$k_{fWf}$	Fracture conductivity, md-ft
$m$	Slope
$P$	Pressure, psi
$P_{wf}$	Well-flowing pressure, psi
$q$	Oil flow rate, STB/D
$q_g$	Gas flow rate, MSCF/D
$r_w$	Wellbore radius, ft
$x_f$	Half-fracture length, ft
$s$	Skin factor. Laplace parameter
$t$	Test time, hr

$T$	Reservoir temperature, °R
$(t*\Delta P')$	Pressure derivative, psi
$(t_D*P_D')$	Dimensionless pressure derivative

## Greek

$\Delta$	Change
$\phi$	Porosity, fraction
$\lambda$	Matrix-fracture network interporosity flow parameter
$\lambda_f$	Fracture network-Hydraulic fracture interporosity flow parameter
$\mu$	Viscosity, cp
$\xi$	Variable to identify homogeneous ( $\xi=1$ ) or heterogeneous ( $\xi=\omega$ ) reservoirs
$\omega$	Dimensionless storativity coefficient

## Suffixes

$2BR$	Birradial flow in the homogeneous region (second birradial)
$2BR$	Second birradial flow at the time of 1 hr
$BR$	Birradial
$BR1$	Birradial at 1 hr
$usBRi$	Intersect of birradial and unit-slope lines
$D$	Dimensionless
$Dxf$	Dimensionless based on half-fractured length
$ell$	Elliptical (same as birradial)
$L1$	Linear at 1 hr
$usLi$	Intersect of linear and unit-slope lines
$max$	Maximum point
$min$	Minimum point
$w$	Well
$t$	Time
$us$	Unit-slope during heterogeneous to homogeneous transition period
$usi$	Intersect of radial al unit-slope lines

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**Appendix-A.** Estimating half-fracture length by Conventional Analysis

Escobar et al. (2014) provided the following equation;

$$\Delta P = 9.4286 \frac{q\mu B}{kh} \left( \frac{k}{\omega\phi\mu c_t x_f^2} \right)^{0.36} t^{0.36} \quad (\text{A.1})$$

Or,

$$\Delta P = m_{ell} t^{0.36} \quad (\text{A.2})$$

Which implies that a Cartesian plot of  $\Delta P$  vs.  $t^{0.36}$  (for drawdown) or  $\Delta P$  vs.  $[(t_p + \Delta t)^{0.36} - \Delta t^{0.36}]$  (for buildup) provides a straight line which slope,  $m_{ell}$ , provides the half-fracture length,

$$x_f = \sqrt{9.4286 \frac{q\mu B}{kh m_{ell}} \left( \frac{k}{\omega\phi\mu c_t} \right)^{0.36}} \quad (\text{A.3})$$

During the homogeneous part, Equation (A.3) becomes:

$$x_f = \sqrt{9.4286 \frac{q\mu B}{kh m_{ell}} \left( \frac{k}{\phi\mu c_t} \right)^{0.36}} \quad (\text{A.4})$$

If the early flow regime is linear, then:

$$\Delta P_{L1} = 4.064 \left( \frac{qB}{h} \right) \sqrt{\frac{\mu t}{\phi c_t k \omega x_f^2}} \quad (\text{A.5})$$

$$\Delta P_{L1} = m_L \sqrt{t} \quad (\text{A.6})$$

So from the Cartesian plot of either pressure or pressure drop versus the square root of time:

$$x_f = 4.064 \left( \frac{qB}{hm_L} \right) \sqrt{\frac{\mu}{\phi c_t k \omega}} \quad (\text{A.7})$$

**Appendix-B. Gas flow**

Equations (1), (2) and (3) are applied to gas wells if the viscosity and total system compressibilities are given at initial conditions, it means  $(\phi c_t)_i$ . However, if those expressions are expressed using the pseudotime concept, Agarwal (1979), it yields:



$$t_{Daxf} = \frac{0.000263kt}{\phi x_f^2} t_a(P) \quad (B.1)$$

For gas wells, Agarwal (1979) also included the pseudopressure definition,

$$m(P)_D = \frac{hk(m(P_i) - m(P))}{1422.52q_{sc}T} \quad (B.2)$$

which dimensionless pseudopressure derivative is given by:

$$t_D * m(P)_D' = \frac{hk(t^* \Delta m(P)')}{1422.52q_{sc}T} \quad (B.3)$$

After replacing Equations (B.1), (B.2) and (B.3) into Equations (3) and (4) leads to obtain:

$$x_f = 529.16 \left( \frac{q_{sc}T}{h(\Delta m(P)_{BR})} \right)^{1.3889} \sqrt{\frac{t_a(P)_{BR}}{\omega \phi k^{1.778}}} \quad (B.4)$$

$$x_f = 138.687 \left( \frac{q_{sc}T}{h(t^* \Delta m(P)'_{BR})} \right)^{1.3889} \sqrt{\frac{t_a(P)_{BR}}{\omega \phi k^{1.778}}} \quad (B.5)$$

$$x_f = 138.687 \left( \frac{q_{sc}T}{h(t^* \Delta m(P)'_{BR1})} \right)^{1.3889} \sqrt{\frac{1}{\omega \phi k^{1.778}}} \quad (B.6)$$

For the homogeneous zone:

$$x_f = 529.16 \left( \frac{q_{sc}T}{h(\Delta m(P)_{2BR})} \right)^{1.3889} \sqrt{\frac{t_a(P)_{2BR}}{\phi k^{1.778}}} \quad (B.7)$$

$$x_f = 138.687 \left( \frac{q_{sc}T}{h(t^* \Delta m(P)'_{2BR})} \right)^{1.3889} \sqrt{\frac{t_a(P)_{2BR}}{\phi k^{1.778}}} \quad (B.8)$$

$$x_f = 138.687 \left( \frac{q_{sc}T}{h(t^* \Delta m(P)'_{2BR1})} \right)^{1.3889} \sqrt{\frac{1}{\phi k^{1.778}}} \quad (B.9)$$

In conventional analysis for gas flow case, Equation (3) becomes,

$$\Delta P = 94.989 \frac{q_{sc}T}{kh} \left( \frac{k}{\omega \phi x_f^2} \right)^{0.36} t_a(P)^{0.36} \quad (B.10)$$

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Then, the slope of the Cartesian plot will give,

$$x_f = \sqrt{94.989 \frac{q_{sc}T}{khm_{ell}} \left( \frac{k}{\omega \phi} \right)^{0.36}} \quad (B.11)$$

For gas wells, Equations (13) and (14) become:

$$x_f = \frac{40.914 q_{sc}T}{h(\Delta m(P))_{L1}} \left( \frac{1}{\omega \phi k} \right)^{0.5} \quad (B.12)$$

$$x_f = \frac{20.457 q_{sc}T}{h(t^* \Delta m(P))'_{L1}} \left( \frac{1}{\omega \phi k} \right)^{0.5} \quad (B.13)$$

The analogous form of Equation (A.7) for gas flow is then:

$$x_f = \frac{40.914 q_{sc}T}{hm_L} \left( \frac{1}{\omega \phi k} \right)^{0.5} \quad (B.14)$$

The gas version of Equations (20) to (23) is:

$$\lambda_f \approx 6560125.65 \left[ \frac{q_{sc}T}{kh[t^* \Delta m(P)'_{\max}] } \right]^{25/9} \quad (B.15)$$

$$\omega = \frac{\lambda_f kt_a(P)_{\max}}{1801.29 \phi x_f^2} \quad (B.16)$$

$$\lambda_f = \frac{1801.29 \omega \phi x_f^2}{kt_a(P)_{\max}} \quad (B.17)$$

Equation (24) for gas flow is:

$$\lambda_f = 20.463 \left( \frac{\phi x_f^2 h [t^* \Delta m(P)]_{us}}{q_{sc} B t_a(P)_{us}} \right)^{25/16} \quad (B.18)$$

Expressions (26) and (27) for gas wells:

$$\omega = 0.0002 \left( \frac{t_a(P)_{\min}}{t_a(P)_{\max}} \right)^2 - 0.0155 \left( \frac{t_a(P)_{\min}}{t_a(P)_{\max}} \right) + 0.2487 \quad (B.19)$$

Dividing Equation (B.6) by Equation (B.9) and solving for the storativity ratio:

$$\omega = \left[ \frac{(t^* \Delta m(P))'_{2BR1}}{(t^* \Delta m(P))'_{BR1}} \right]^{2.7778} \quad (B.20)$$



Dividing Equation (B.13) by Equation (B.9) and solving for the storativity ratio:

$$\omega = 0.0218 \left( \frac{h}{q_{sc} B} \right)^{0.778} \left[ \frac{[t^* \Delta m(P)]_{2.BR1}^{1.3889}}{[t^* \Delta m(P)]_{L1}} \right]^2 k^{0.778} \quad (B.21)$$

Equations (29) and (31) rewritten for gas flow are:

$$\omega = \frac{1}{1456.69} \left( \frac{k \lambda_f t_a (P)_{usBRi}}{\phi x_f^2} \right)^{9/16} \quad (B.22)$$

$$\omega = \frac{6356.314 \phi x_f^2}{k \lambda_f^{32/25} t_a (P)_{usLi}} \quad (B.23)$$

Finally, the gas versions for Equation (33) and (34) are:

$$\lambda_f = 345195361.5 \left[ \frac{q_{sc} T}{kh \Delta m(P)_{pss}} \right]^{25/9} \quad (B.24)$$

$$\omega = 0.3625 e^{-0.9877 \frac{\Delta m(P)_{high}}{\Delta m(P)_{low}}} \quad (B.25)$$